

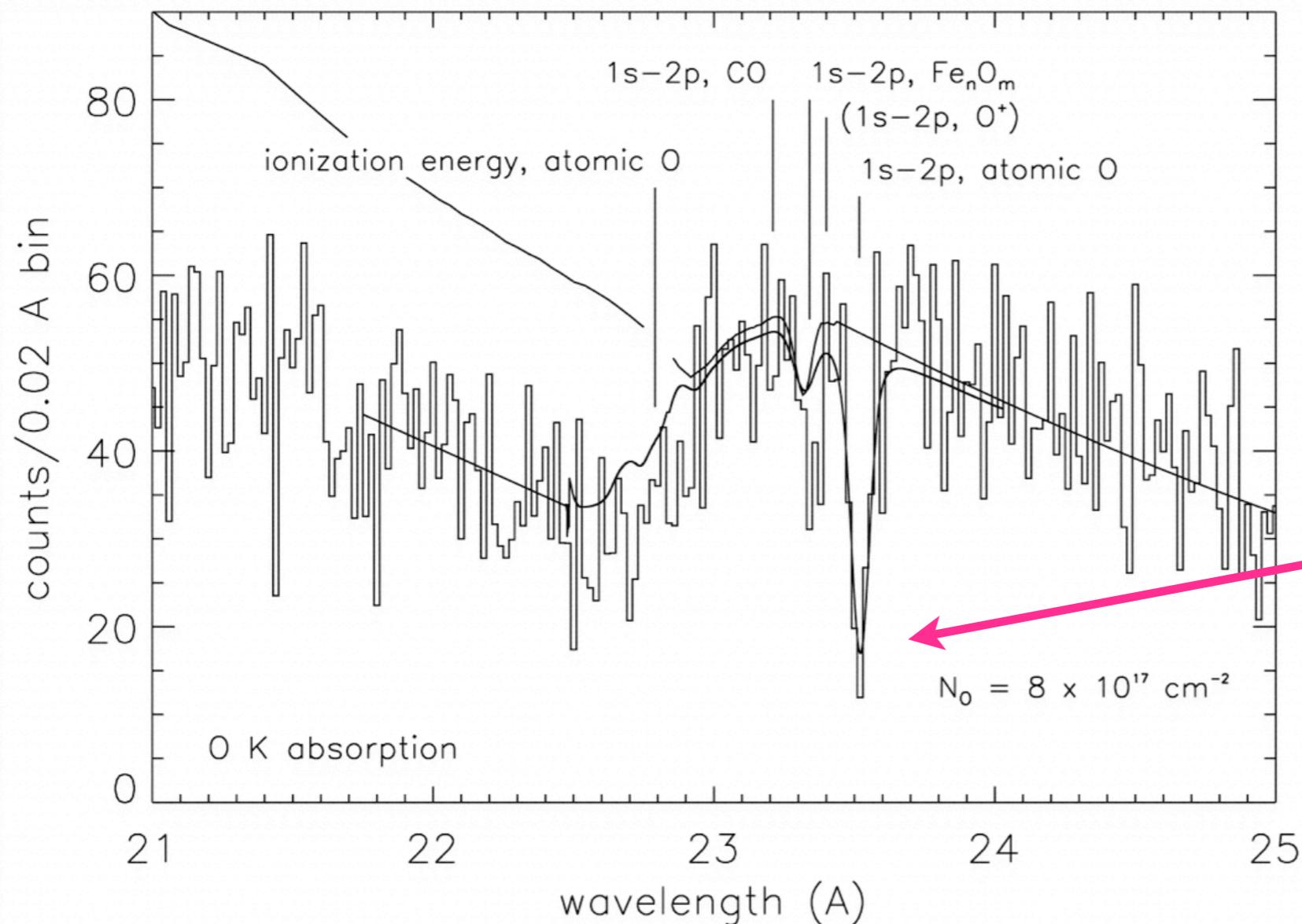
# Absorption Spectroscopy

Frits Paerels (\*)/Columbia University  
AtomDB Workshop, Tokyo Metropolitan University  
September 2014

(\*) standing on the shoulders of many others



Anything that changes the structure of the valence shell (ionization; chemical binding) shifts its energy by  $\Delta E \sim \text{eV per electron}$ : position of absorption lines changes by  $\sim 1$  part in few 100 - 1000



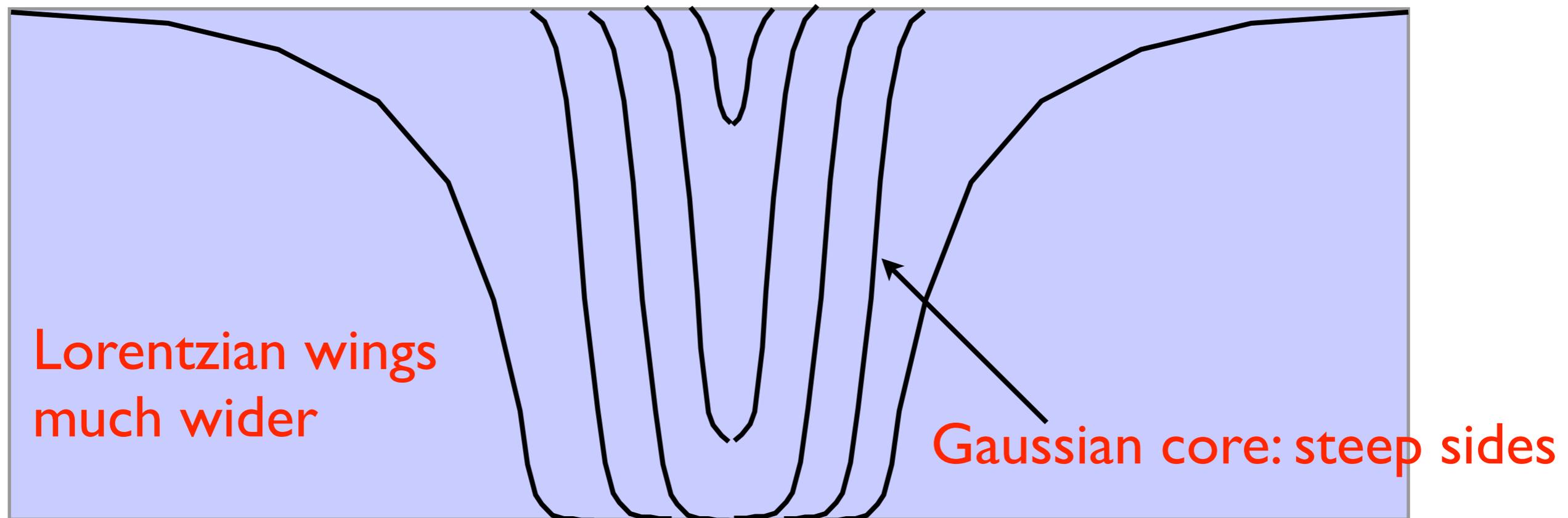
X0614+091  
Chandra LETGS  
FP et al. 2001

saturated!  
( $N_{\text{O}}$  derived  
from edge depth)

absorption lines more sensitive than continuum edges  
to small absorber column densities

convert absorption line 'strength' into column density:  
straightforward, but be careful with saturation

absorption line saturation



absorption line 'strength': *equivalent width (EW)*  
(area of equivalent rectangle under continuum)  
*(NB: independent of spectrometer resolution!!!)*

Saturation EW  $\sim$  line core width;  
often (??) thermal velocity dispersion:

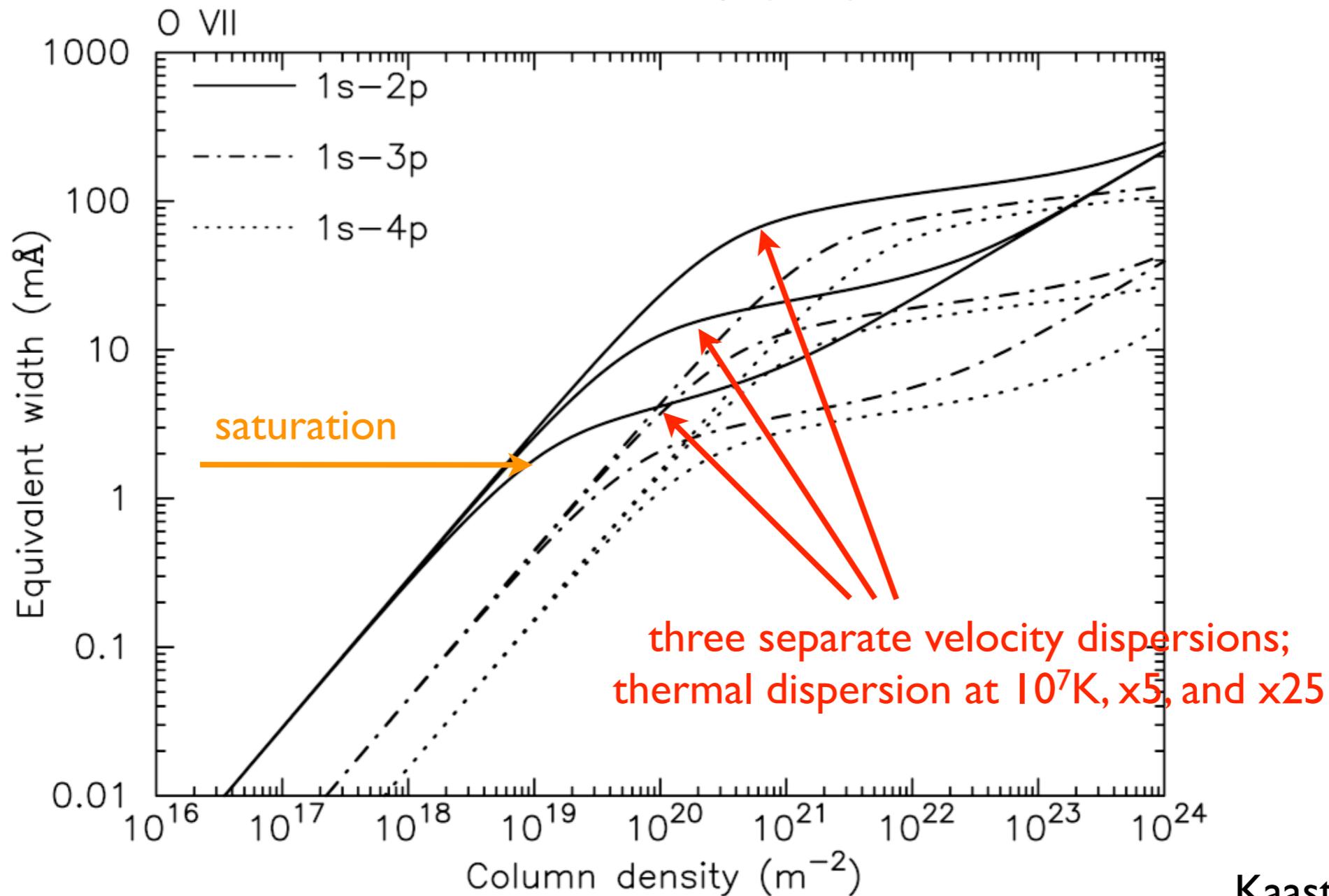
$$\frac{v}{c} \sim \left( \frac{kT}{M_i c^2} \right)^{1/2} = \frac{1}{13000} \left( \frac{M_i}{16m_p} \right)^{-1/2} \left( \frac{T}{10^6 \text{ K}} \right)^{1/2}$$

moderate saturation basically not diagnosable in  
*unresolved* lines

circumvent: see higher order series members: do 'COG'

# the 'Curve of Growth'

'curve of growth': EW vs. absorber column density

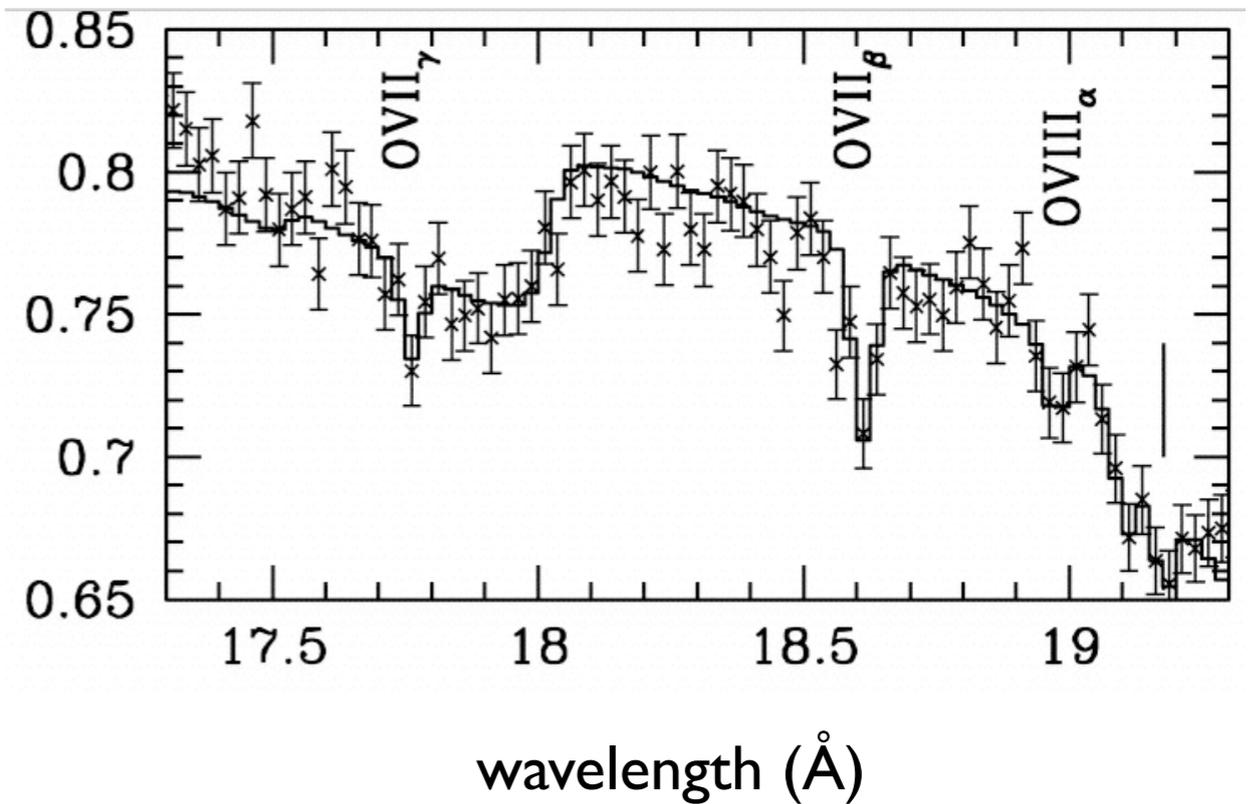
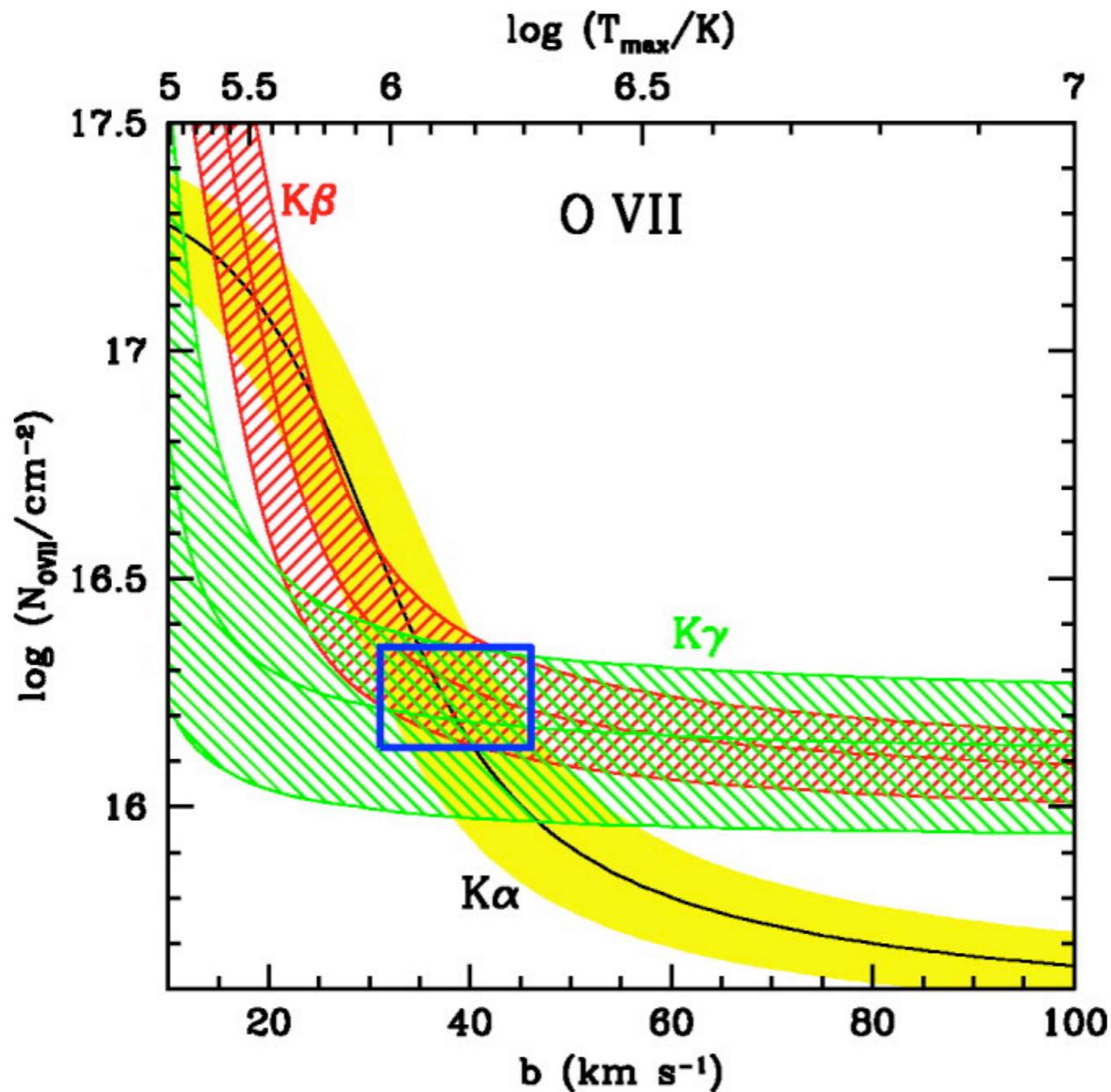


Kaastra et al. 2008

**NB: resolving 1 mA in O VII resonance line requires  $R \sim 20,000!!!$**

higher order series members:

cross section typically order of magnitude smaller than  $n=1-2$ :  
saturate at much higher column density



circumgalactic absorption, Mkn 421 *Chandra* LETGS; Williams *et al.* 2005

# Photoelectric Absorption Edges

1. Far above the ionization energy
2. Close to the ionization energy; XANES, (E)XAFS, ...

1. Far above the ionization energy  
(but not relativistic)

perturbation theory:

$$\text{cross section} \propto \left| \langle \psi_{\text{final}} | H_{\text{int}} | \psi_{\text{initial}} \rangle \right|^2$$

free photoelectron in  
Coulomb field of ion

bound atomic state

incident radiation field

free electron  $\sim$  plane wave;

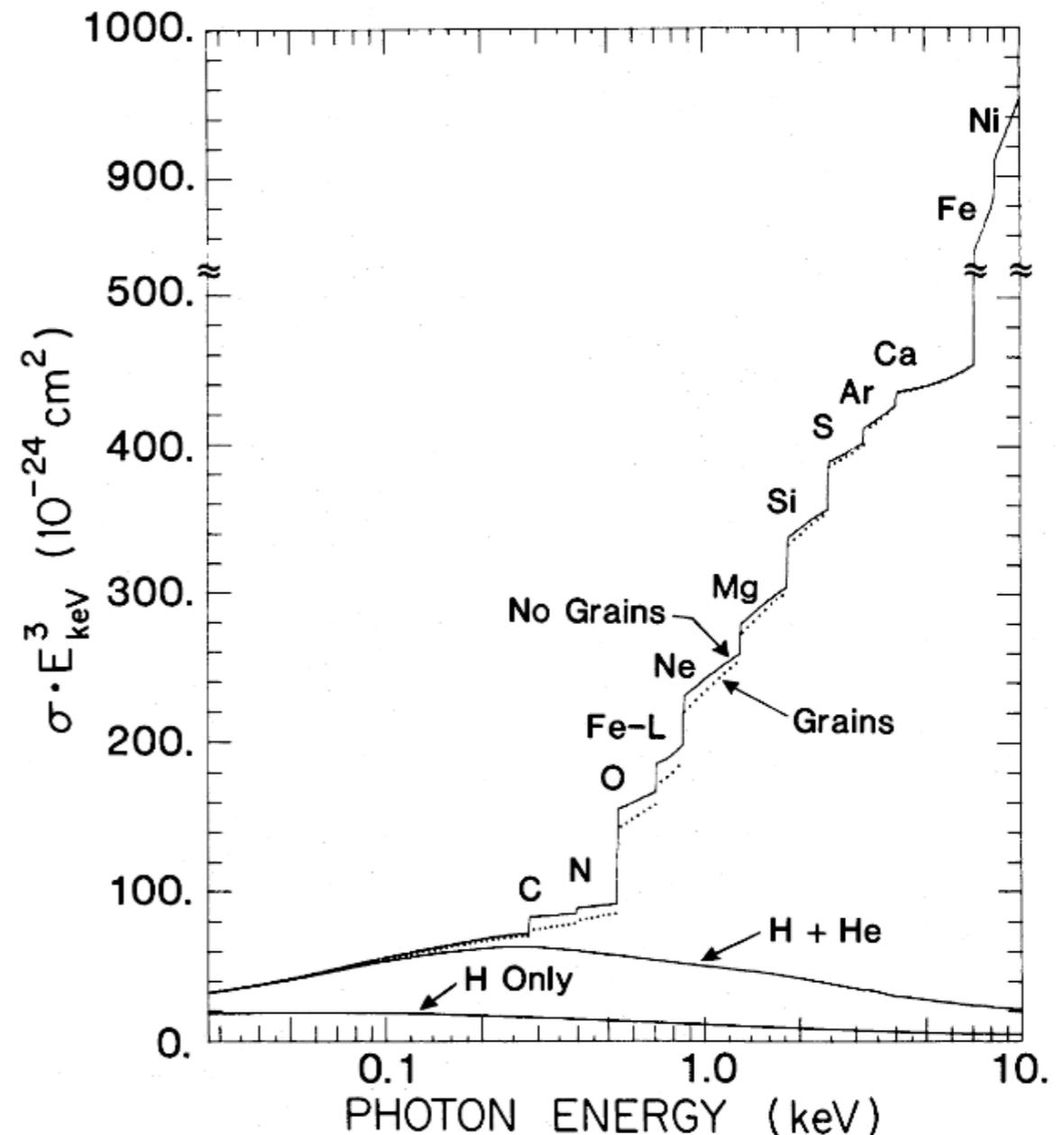
dipole approximation for radiation field:  $\exp(i\mathbf{k}\cdot\mathbf{r}) \rightarrow 1$

$$\sigma = 4\sqrt{2} \sigma_T \alpha^4 Z^5 \left( \frac{m_e c^2}{\hbar\omega} \right)^{7/2}$$

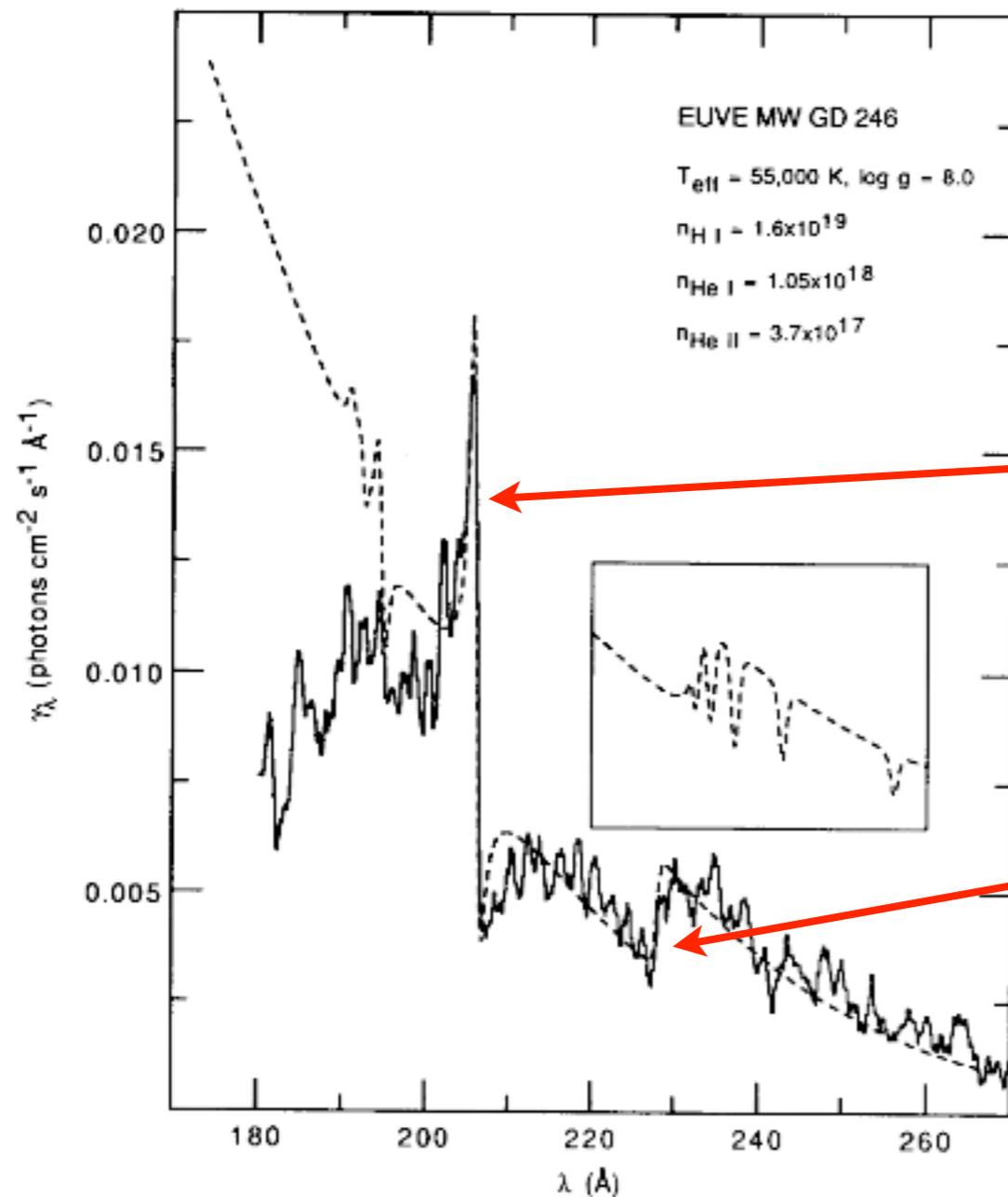
Thomson cross section

at K threshold:  $E_K = Z^2 \cdot Ry$ , so  
 $\sigma \sim \text{constant } Z^{-2}$

cross section of the ISM;  
Morrison & McCammon 1983



# Ia. Intermezzo: autoionizing resonances in continuum cross sections



interstellar absorption in hot DA white dwarf GD246

He I autoionization resonance (ISM)

He II Ly edge (ISM)

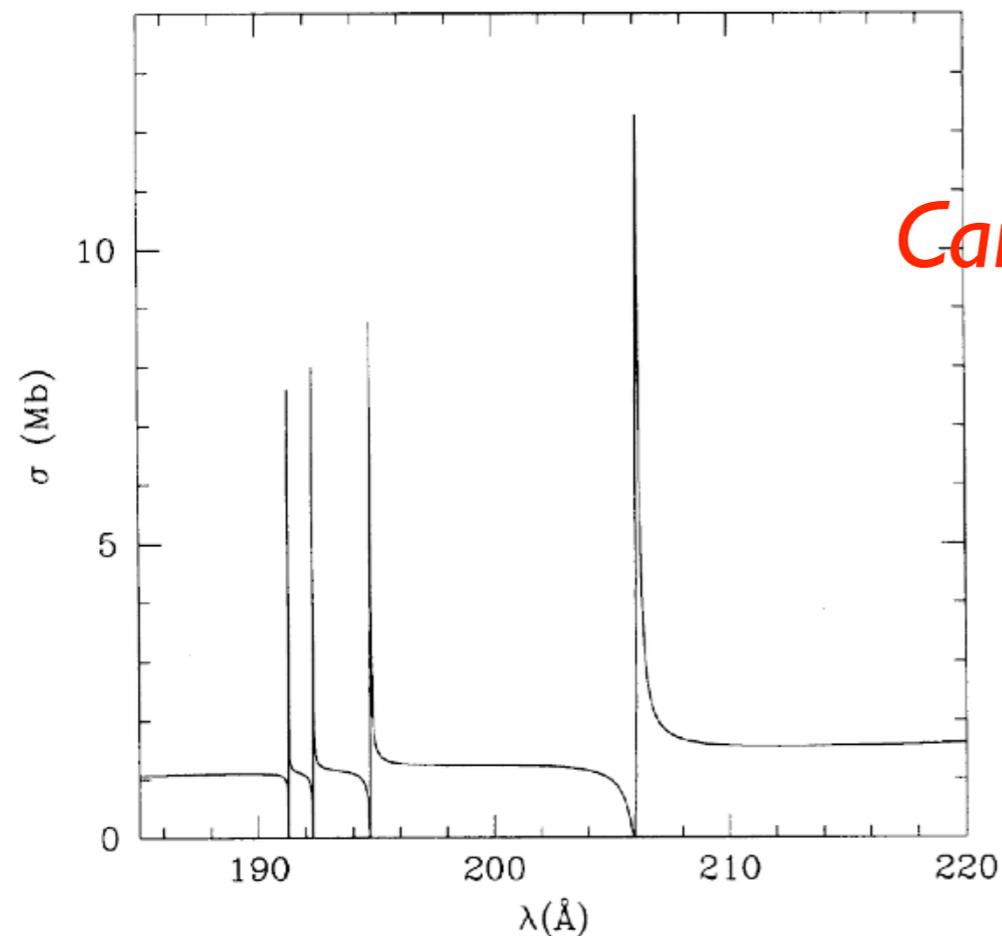
FIG. 1.—Pure H synthetic spectrum adjusted to the observed EUV spectrum of GD 246 in the range  $\lambda \geq 200 \text{ \AA}$  (see § 3.1). *Inset:* Synthetic spectrum with a trace abundance of helium ( $\text{He}/\text{H} = 2 \times 10^{-5}$ ) illustrating pressure broadening effects near the He II series limit in a model atmosphere.

# autoionizing resonances in continuum cross sections



$2\ell n\ell'$ ; e.g.  $2s2p$ ;  $E < \text{ionization energy}$

autoionization;  $h\nu$  destroyed



*Can we see this in C V, O VII, Ne IX, ...?*

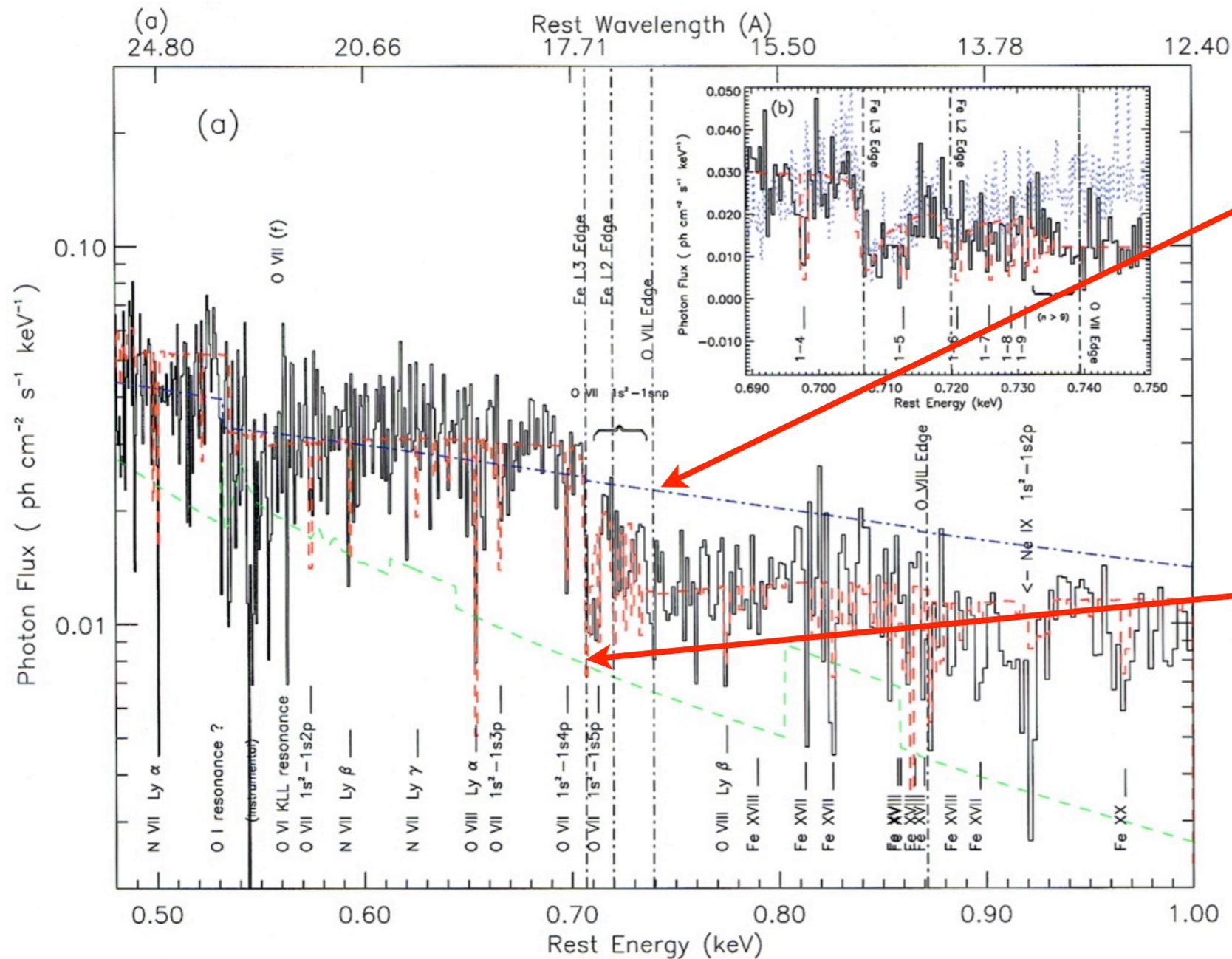
FIG. 1. The photoionization cross section of He I showing the first four  $(sp, 2n+)^1P^0$  autoionizing resonances. Cross sections are given in Mb ( $10^{-18}$  cm).

He I; Rumph *et al.* 1994

# 2a. Close to the ionization energy (atomic absorption)

## Few-electron ions: example O VII

O VII series limit 'eroded' by presence of multiple, weak, unresolved lines converging on limit

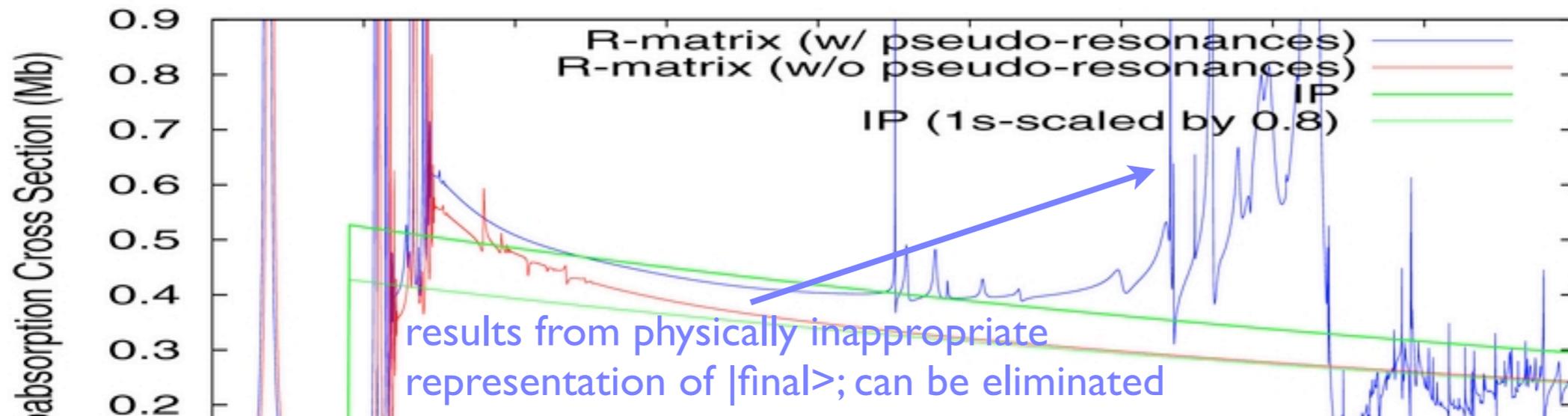
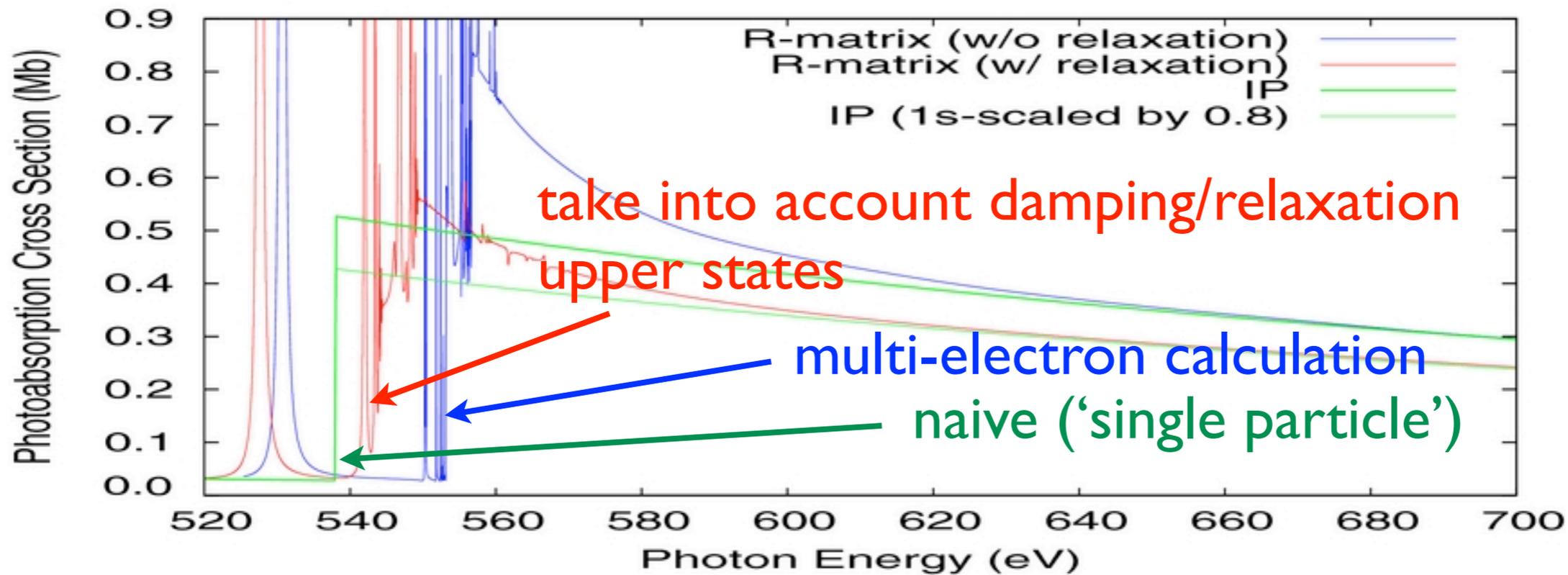


Also note presence of Fe L 2p<sub>1/2</sub>, 2p<sub>3/2</sub> "edges" (see later)

# 2a. Close to the ionization energy (atomic absorption)

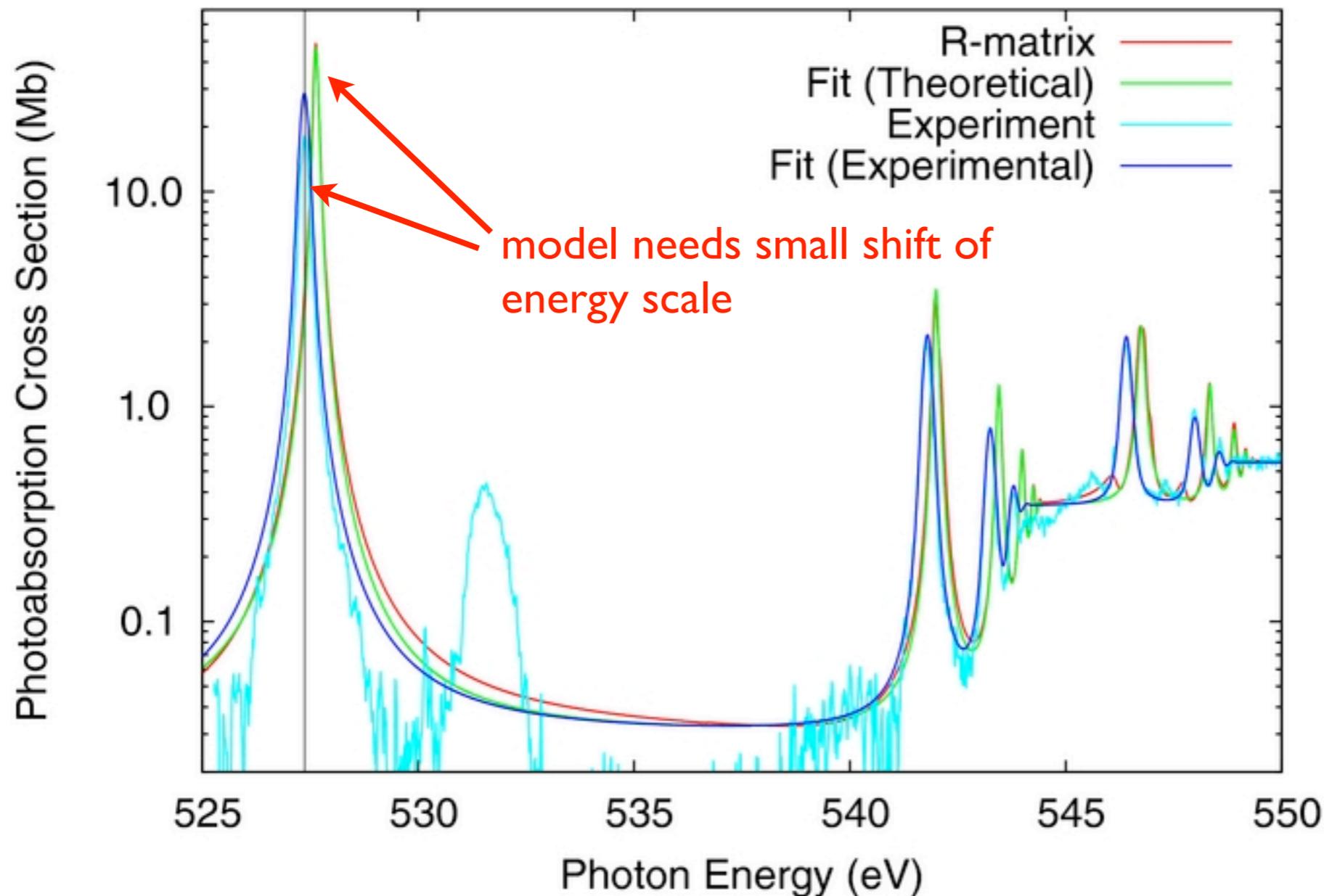
## many-electron ions/atoms

O I K absorption; R-matrix, Gorczyca *et al.* 2013;  
*note importance of detailed treatment of all possible decays from upper state!*  
(blue  $\rightarrow$  red)



## 2a. Close to the ionization energy (atomic absorption) many-electron ions/atoms

zoom in; compare to experiments  
(difficult to make monatomic O!!)

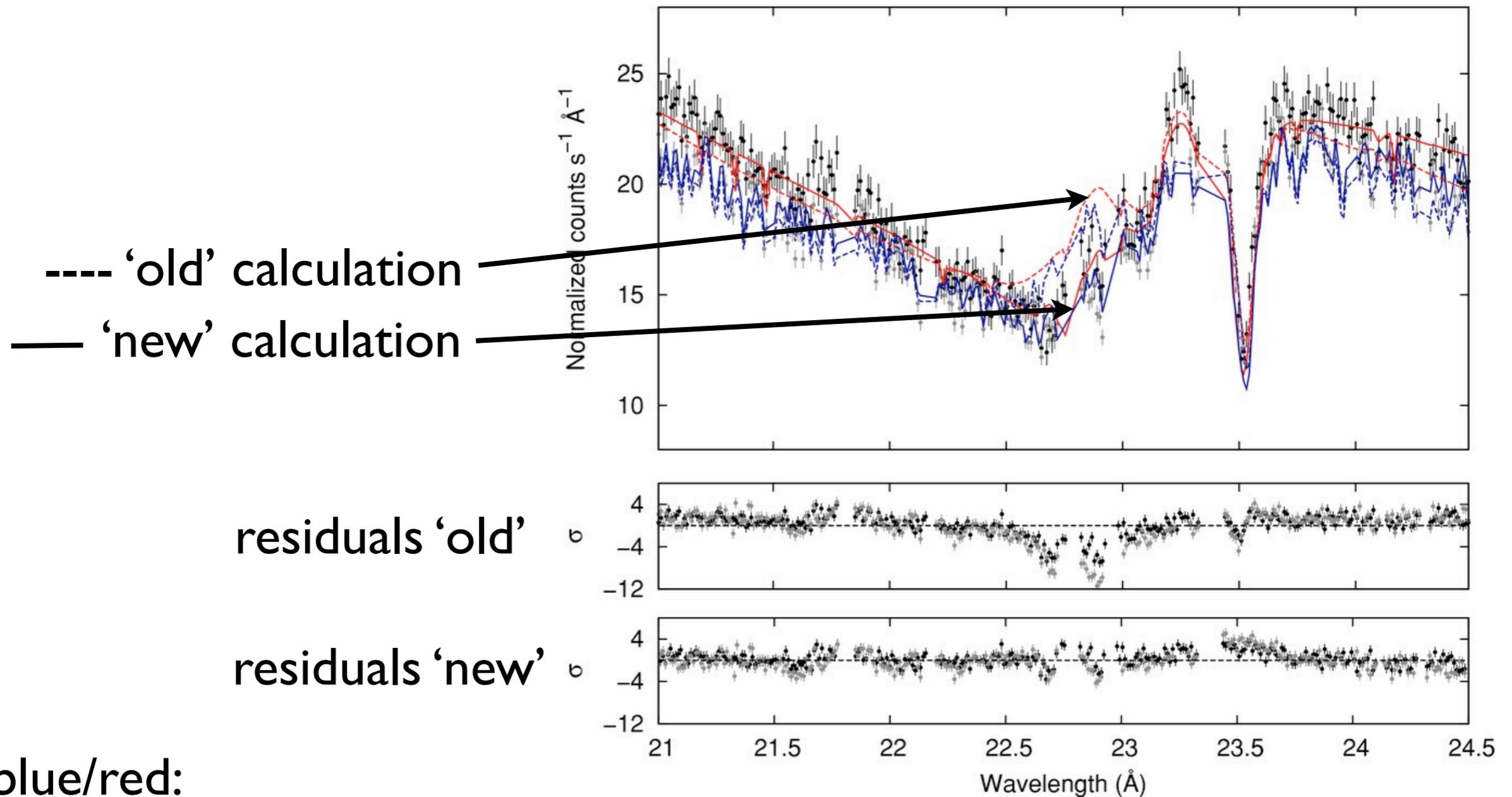


model convolved with  
0.18 eV FWHM Gauss  
to match experimental  
resolution

# 2a. Close to the ionization energy: atomic absorption many-electron ions/atoms

how well does this stuff do?

Sco X-1 XMM/RGS



---- 'old' calculation

— 'new' calculation

residuals 'old'

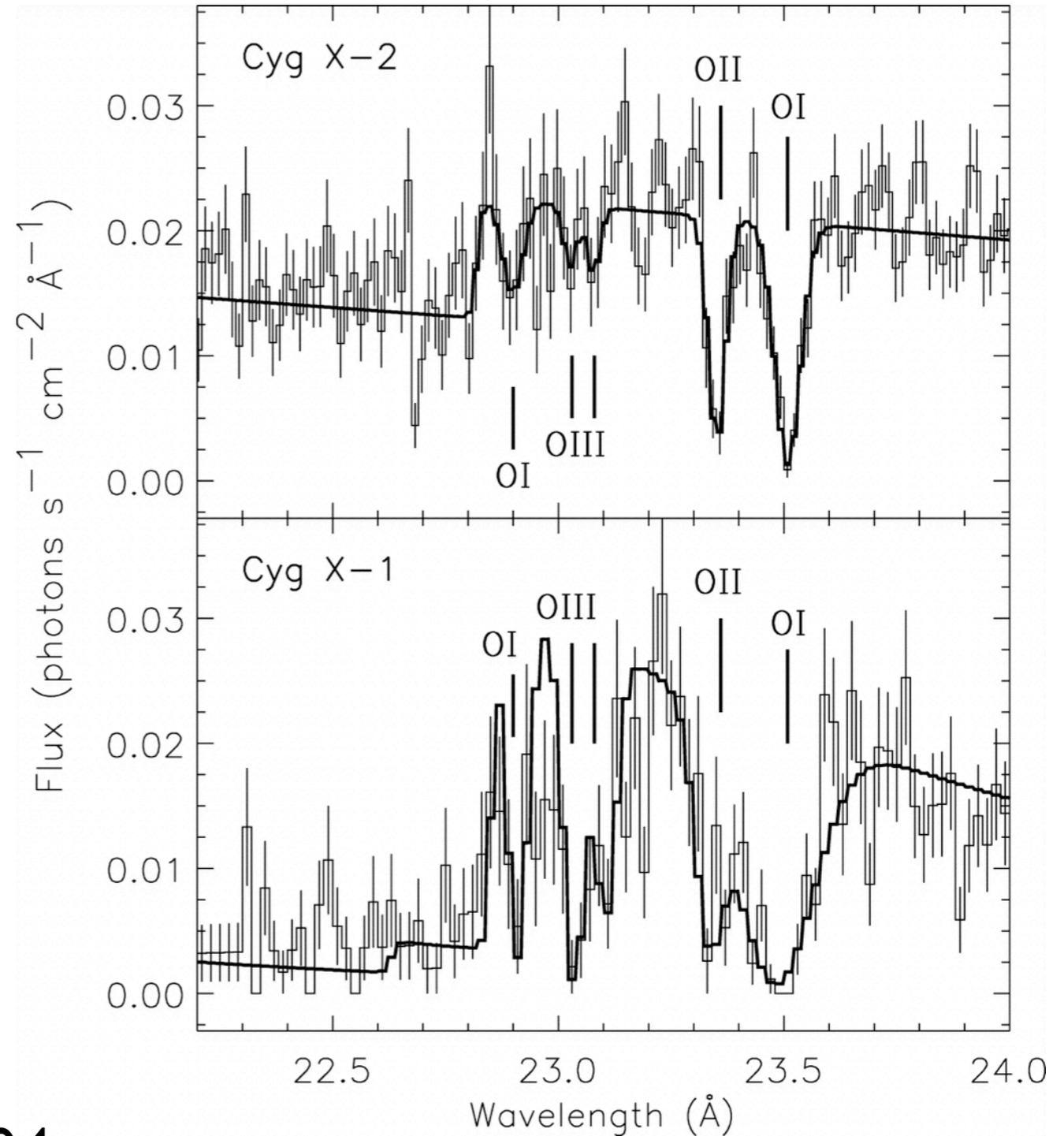
residuals 'new'

blue/red:

two different datasets

# 2a. Close to the ionization energy: atomic absorption many-electron ions/atoms

effect of ionization



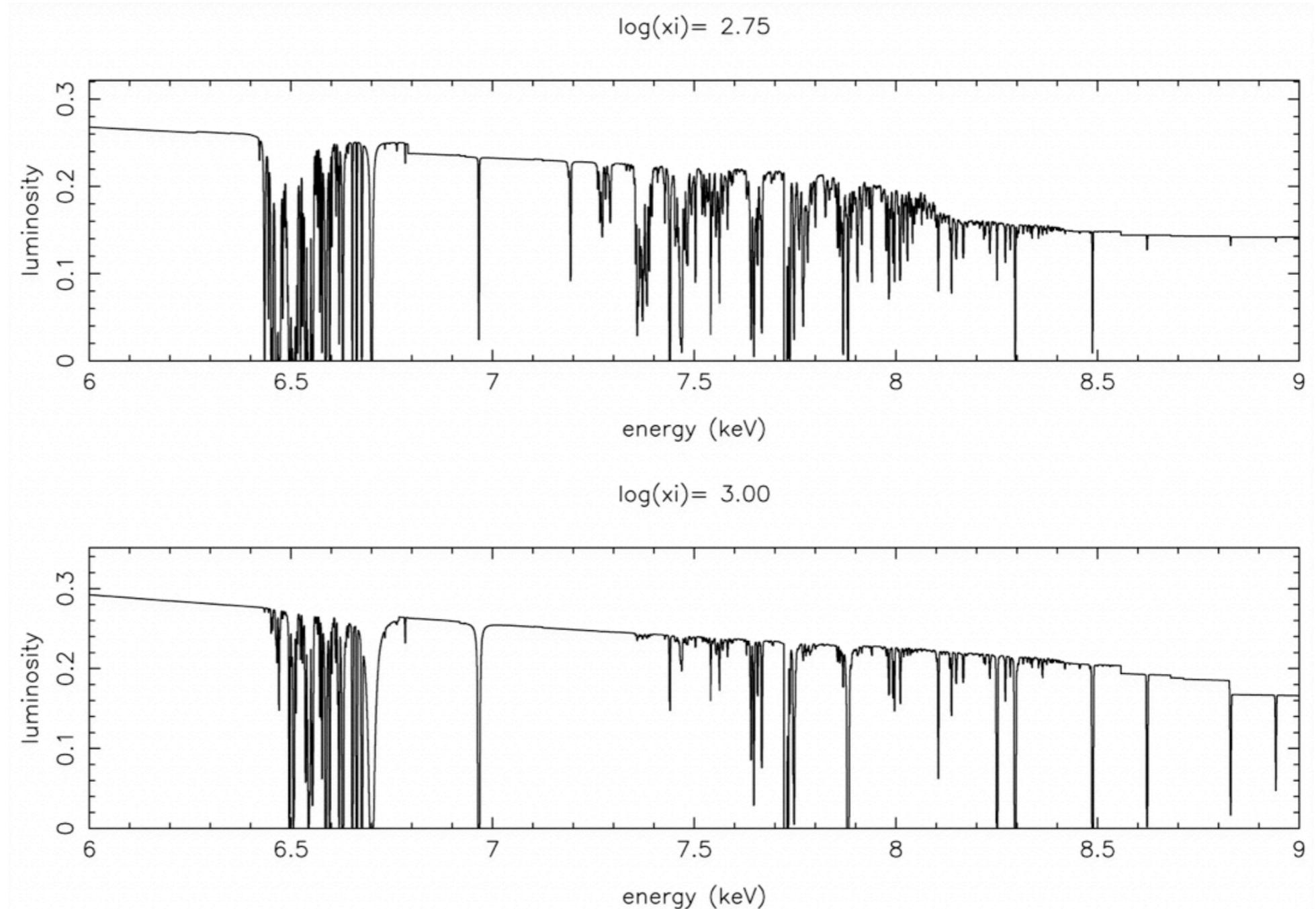
# 2a. Close to the ionization energy: atomic absorption

## many-electron ions/atoms

### innershell absorption by Fe

gas in photo-ionization equilibrium;  
fixed ionization parameter

this example:  
mostly L-shell ions contributing



Kallman *et al.* 2004

2b. Close to the ionization energy:  
absorption in molecular and solid phase  
XANES, (E)XAFS, ...

definition of 'close':

photoelectron momentum:

photoelectron de Broglie-wavelength

interatomic distance

$$\frac{\hbar^2 k^2}{2m_e} = E_\gamma - \chi; \lambda_e \sim a \Rightarrow$$

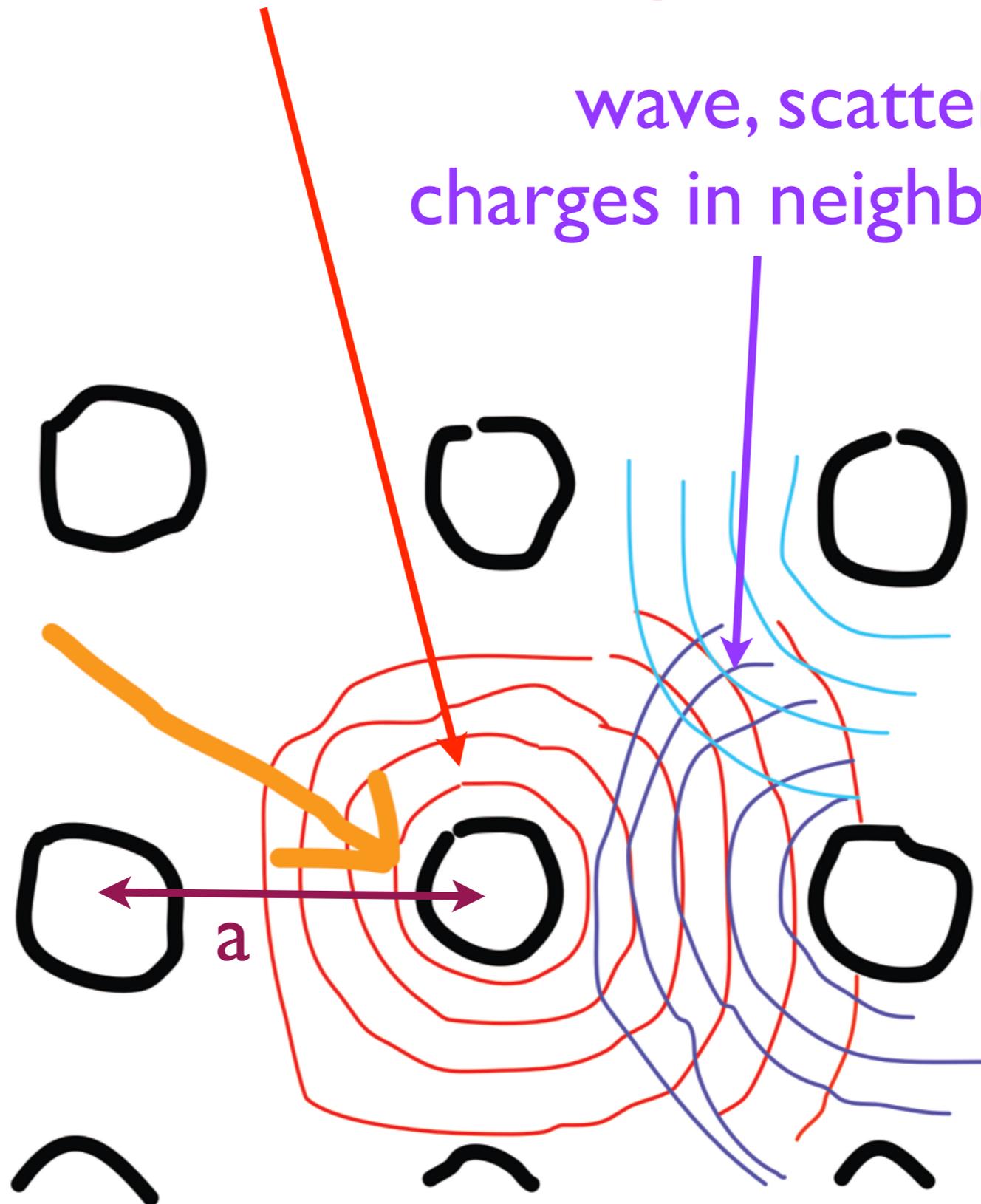
$$\Rightarrow E_\gamma - \chi \sim \frac{h^2}{2m_e a^2} = 150(a/1 \text{ \AA})^{-2} \text{ eV}$$

why does interatomic distance matter?

outgoing photoelectron de Broglie wave

wave, scattered by charges in neighboring atom

photon



# why does interatomic distance matter?

transition probability  $\propto$

$$\propto |\langle \psi_{\text{final}} | H_{\text{int}} | \psi_{\text{initial}} \rangle|^2 = |\langle \psi_{\text{out}} + \psi_{\text{scattered}} | H_{\text{int}} | \psi_{\text{initial}} \rangle|^2$$

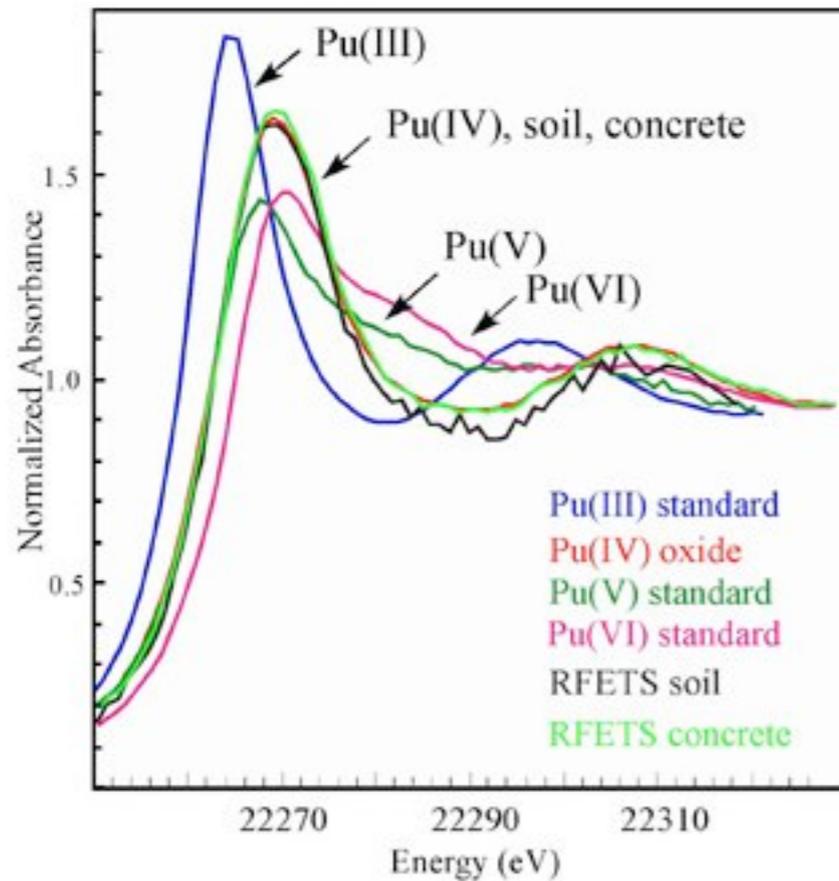
interference! modulates the entire matrix element;  
when amplitude of  $\psi_{\text{out}} + \psi_{\text{scattered}}$  is small near  
absorbing atom (where  $\psi_{\text{initial}}$  is nonzero):  
*transition probability suppressed*

‘resonance’:  $n\lambda_e \sim a$

note: this does *not* require a full lattice! only  
short-range order!

Also:  $\psi_{\text{scattered}}$  is sensitive to nature of scatterer!

from Wikipedia



absorption spectrum of Pu  
from samples taken near Rocky Flats

the effect occurs not just in solids, but in any  
medium with short-range order  
(such as molecules in gas phase!)

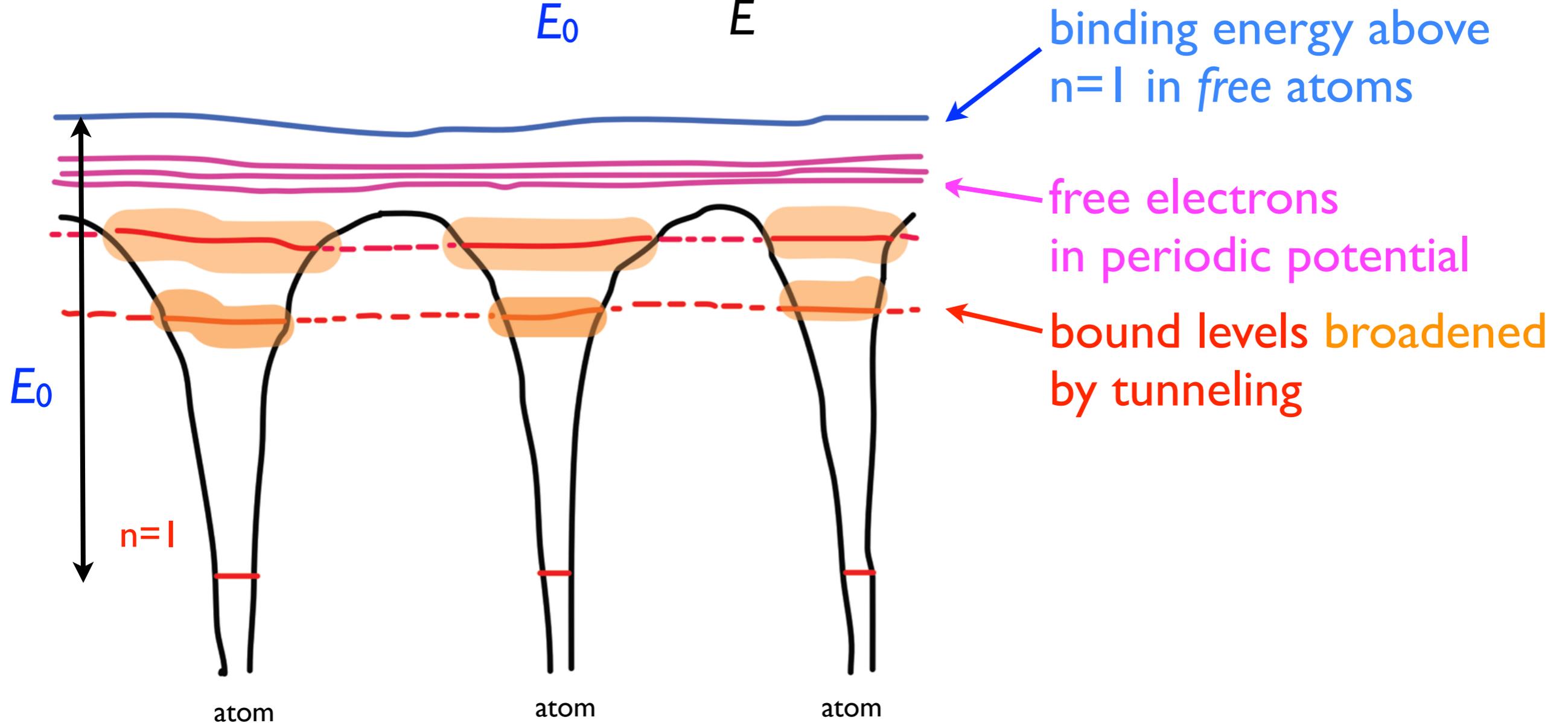
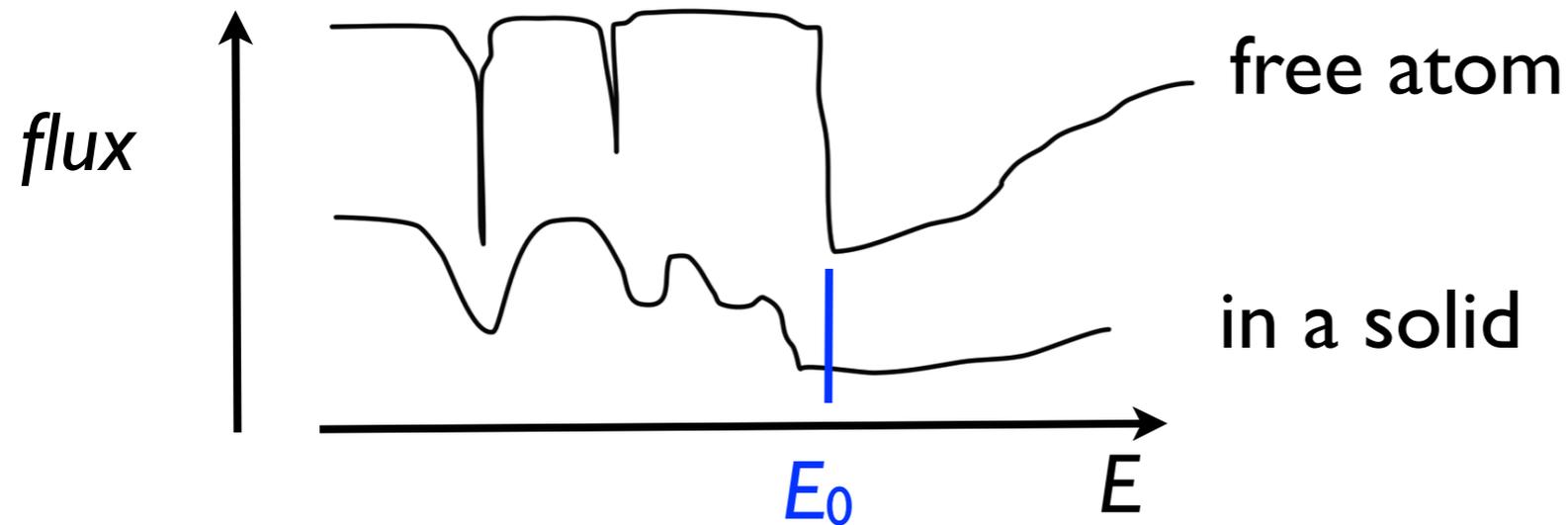
terminology:

XANES: X-ray absorption Near-Edge Structure

(E)XAFS: (Extended) X-ray Absorption Fine Structure

if  $\lambda_e \gtrsim a$ : XANES; if  $\lambda_e < a$ : XAFS

# another contributor to XANES: 'band structure'

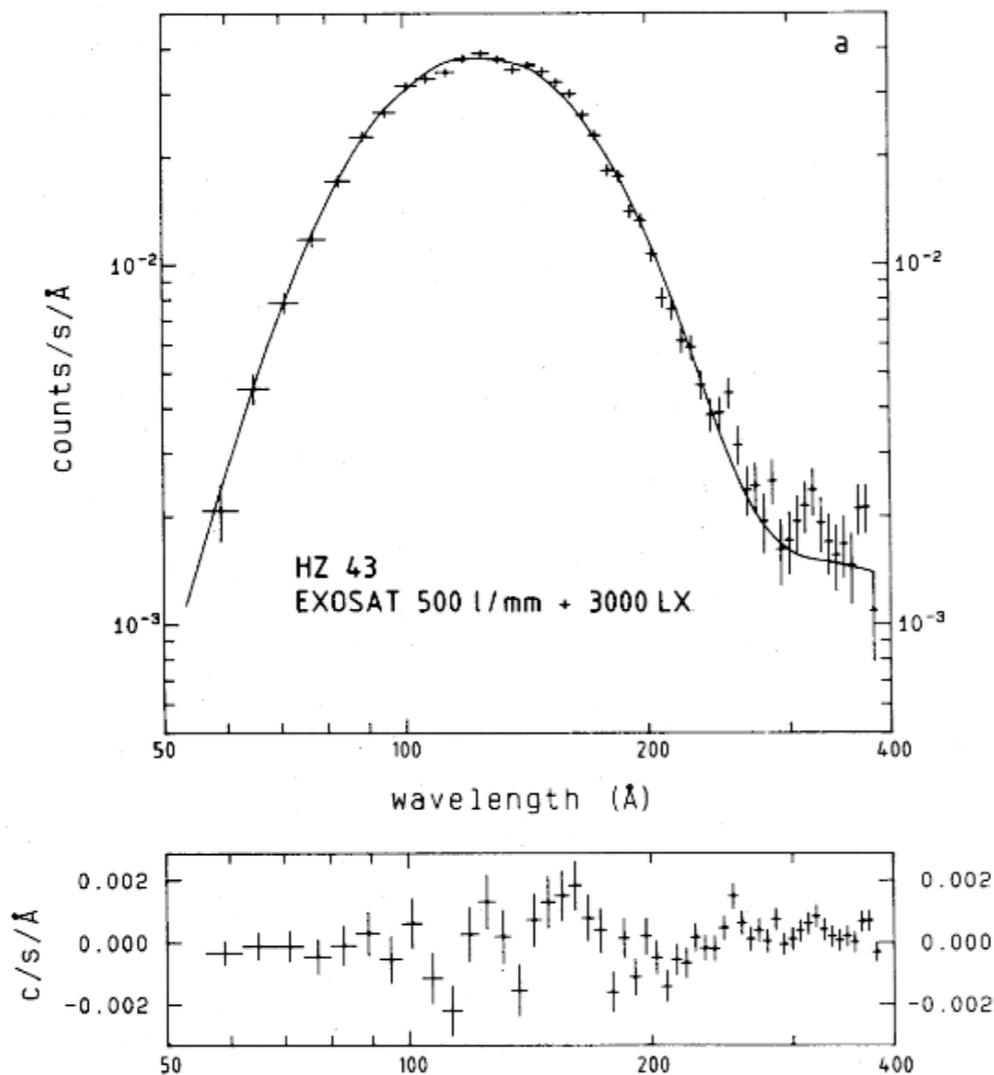


obvious importance to astrophysics:

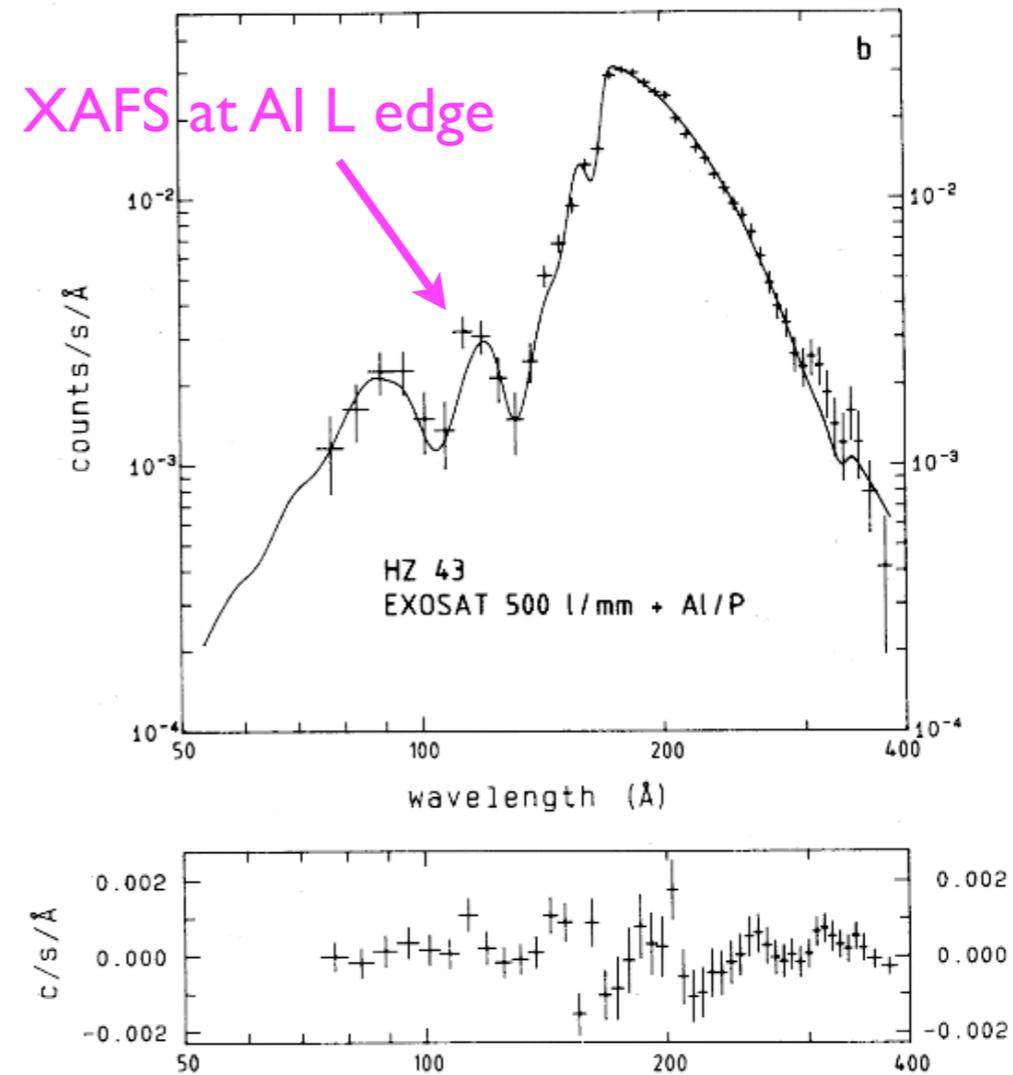
(1) instrumental calibration...

(2) molecular and dust absorption in ISM

plastic beam filter



Al beam filter



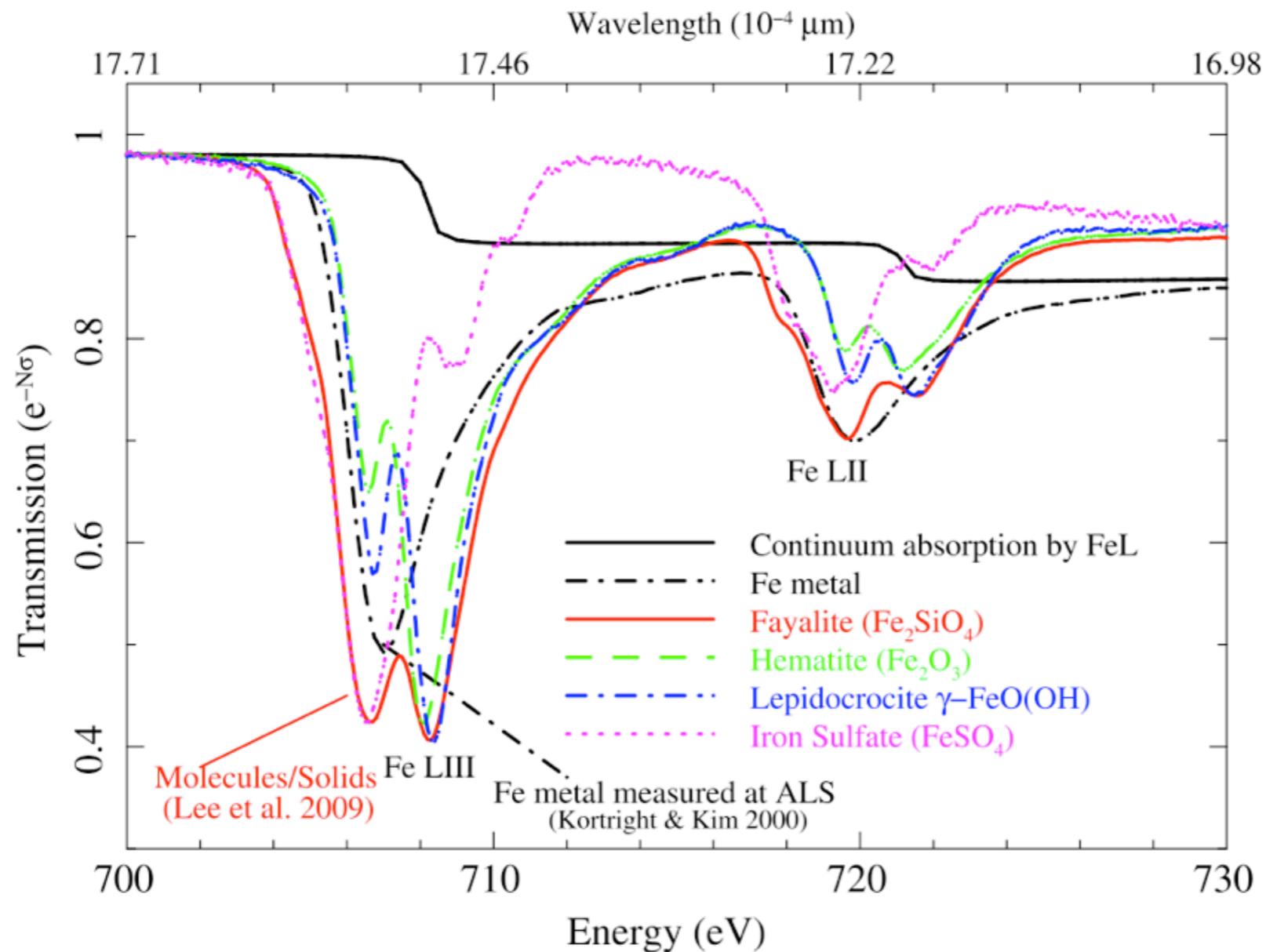
(very old) example of an instrument calibration issue

hot DA white dwarf HZ43, EXOSAT TGS

FP et al. 1986

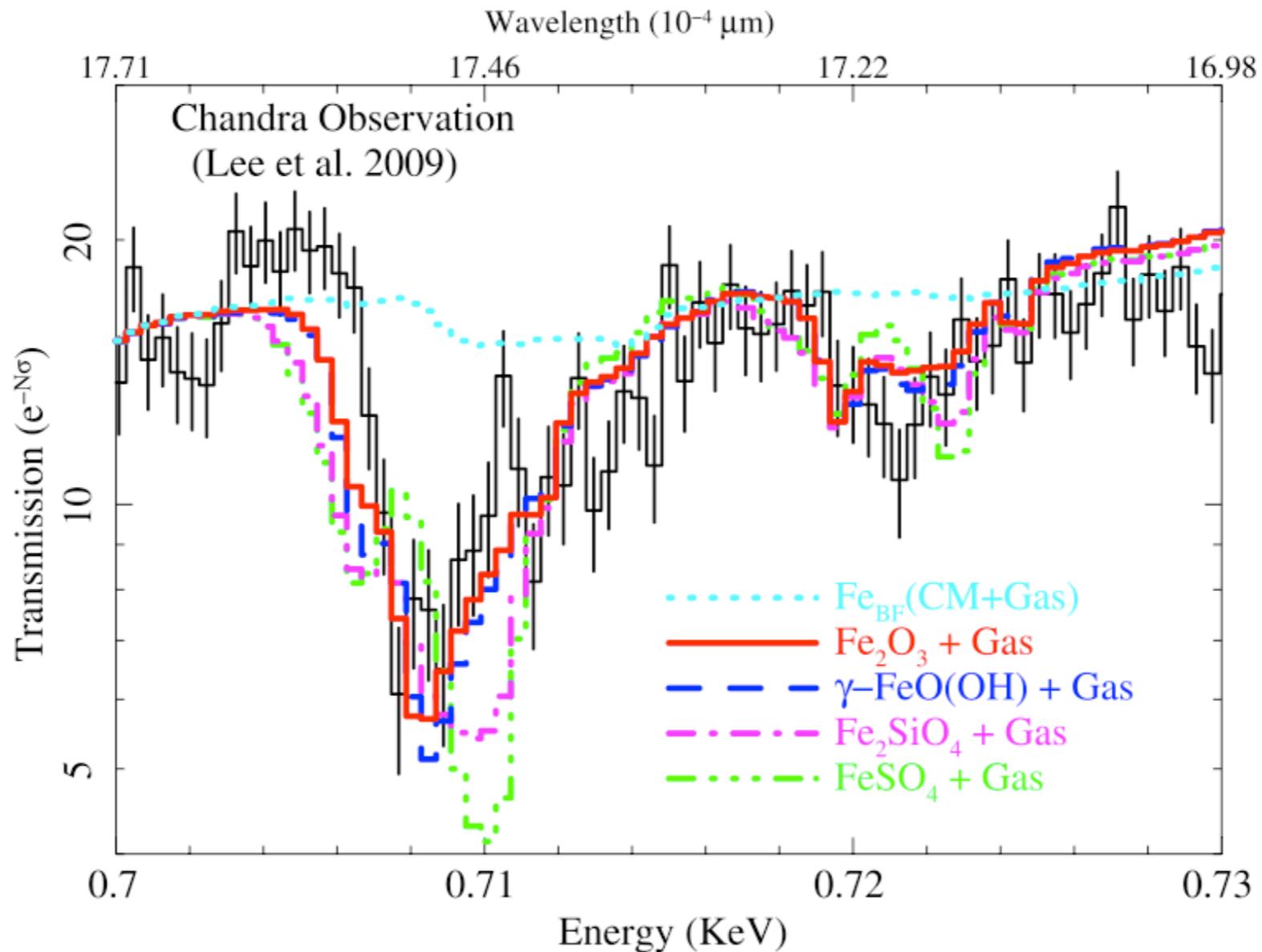
# Interstellar Absorption by molecules and dust

## Fe L absorption: the L 'edges'



**Fig. 1** Absorption calculated from laboratory cross section ( $\sigma$ ) measurements for 0.7 keV FeL presented in Lee et al. (2009). At this high spectral resolution ( $R = 3000$ , the baseline IXO value), different molecules (color), metals (dashed-dot black) and continuum absorption (solid black), are easily distinguished

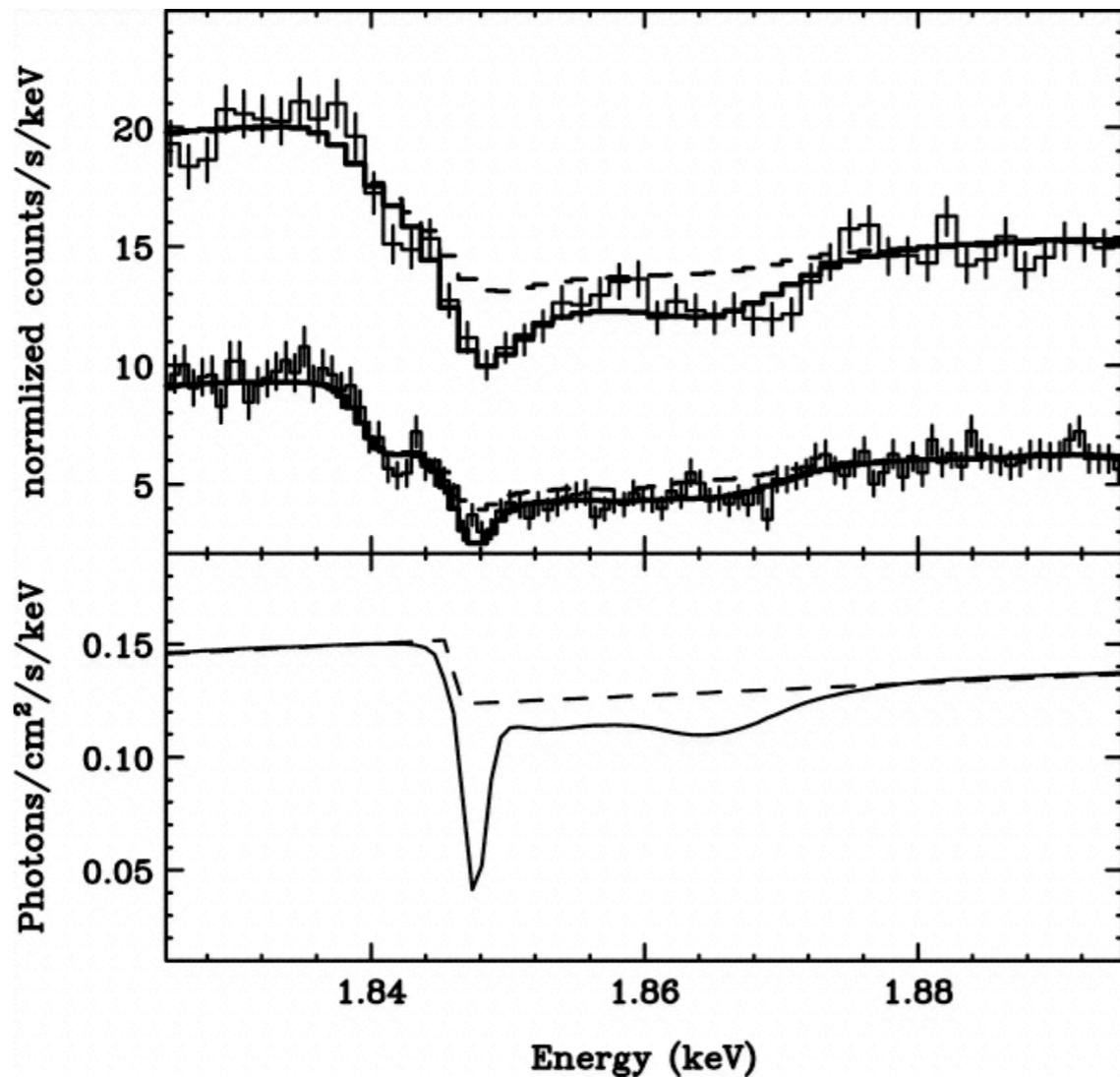
# Interstellar Absorption by molecules and dust



**Fig. 2** A 15ks Chandra HETGS observation of Cygnus X-1 (*black*) zoomed in on the  $\sim 700$  eV ( $\sim 17.7$  Å) FeL spectral region shows that at the current highest available spectral resolution, we can discern at high confidence between Fe<sub>2</sub>O<sub>3</sub> (*red*) and FeSO<sub>4</sub> (*green*), but not between e.g. Fe<sub>2</sub>O<sub>3</sub> and FeO(OH) (*dark blue*). In *light blue* is continuum absorption by Fe L from both gas and solids (CM). In addition, even for a source as bright as Cyg X-1, a 15 ks integration is require to achieve a S/N  $\sim 5$  per bin. Future missions with higher throughput will allow us to engage in such studies with more sources, and accompanying higher spectral resolution as e.g. that shown in Fig. 1 will enable more involved studies of dust structure, not currently possible

# Interstellar Absorption by molecules and dust

real XAFS?



structure at the Si K edge

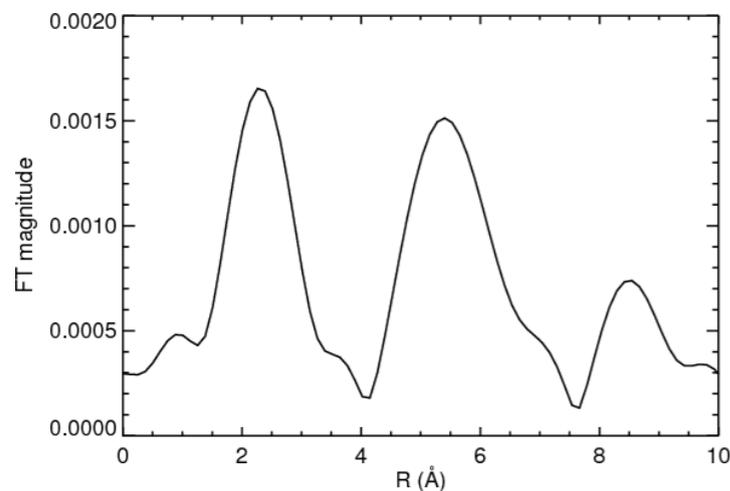
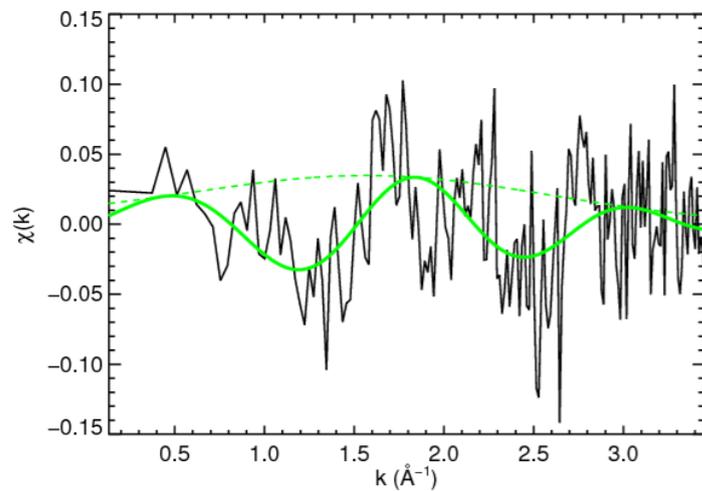
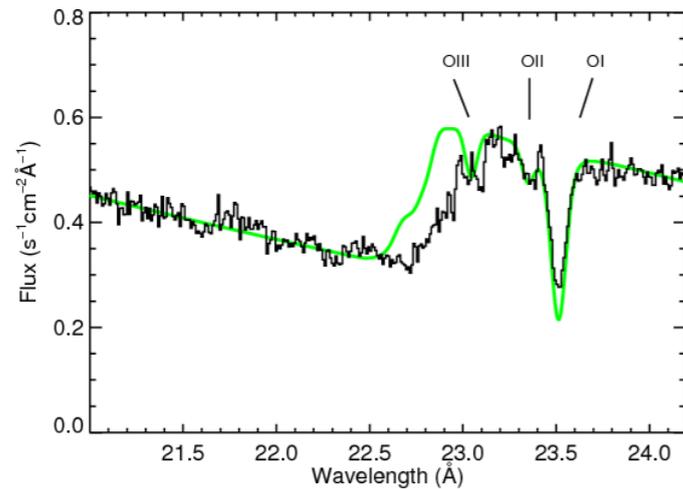
GX 13+1 *Chandra* HETG, METG

Ueda *et al.* 2005

claim: structure not seen in Cyg X-2 (low NH)

# Interstellar Absorption by molecules and dust

real XAFS?



structure at the O K edge  
*exciting (though puzzling), because  
consistent with O in water ice!*  
but signal disappears with use of  
more sophisticated atomic O  
cross section...

Sco X-1 XMM-Newton/RGS  
de Vries & Costantini 2009  
modulation at O K